

Available online at www.sciencedirect.com

ScienceDirect

Linear Algebra and its Applications 420 (2007) 198–207

**LINEAR ALGEBRA
AND ITS
APPLICATIONS**

www.elsevier.com/locate/laa

Linear transformations that are tridiagonal with respect to both eigenbases of a Leonard pair

Kazumasa Nomura ^{a,*}, Paul Terwilliger ^b
^a *College of Liberal Arts and Sciences, Tokyo Medical and Dental University, Kohnodai,
Ichikawa 272-0827, Japan*
^b *Department of Mathematics, University of Wisconsin, 480 Lincoln Drive, Madison, WI 53706, USA*

Received 18 May 2006; accepted 3 July 2006

Available online 22 August 2006

Submitted by R.A. Brualdi

Abstract

Let V denote a vector space with finite positive dimension. We consider a pair of linear transformations $A : V \rightarrow V$ and $A^* : V \rightarrow V$ that satisfy (i) and (ii) below:

- (i) There exists a basis for V with respect to which the matrix representing A is irreducible tridiagonal and the matrix representing A^* is diagonal.
- (ii) There exists a basis for V with respect to which the matrix representing A^* is irreducible tridiagonal and the matrix representing A is diagonal.

We call such a pair a *Leonard pair* on V . Let \mathcal{X} denote the set of linear transformations $X : V \rightarrow V$ such that the matrix representing X with respect to the basis (i) is tridiagonal and the matrix representing X with respect to the basis (ii) is tridiagonal. We show that \mathcal{X} is spanned by

$$I, A, A^*, AA^*, A^*A$$

and these elements form a basis for \mathcal{X} provided the dimension of V is at least 3.

© 2006 Elsevier Inc. All rights reserved.

AMS classification: 05E35; 05E30; 33C45; 33D45

Keywords: Leonard pair; Tridiagonal pair; q -Racah polynomial; Orthogonal polynomial

* Corresponding author.

E-mail addresses: nomura.las@tmd.ac.jp, knomura@pop11.odn.ne.jp (K. Nomura), terwilli@math.wisc.edu (P. Terwilliger).

1. Leonard pairs

We begin by recalling the notion of a Leonard pair. We will use the following terms. A square matrix X is said to be *tridiagonal* whenever each nonzero entry lies on either the diagonal, the subdiagonal, or the superdiagonal. Assume X is tridiagonal. Then X is said to be *irreducible* whenever each entry on the subdiagonal is nonzero and each entry on the superdiagonal is nonzero. We now define a Leonard pair. For the rest of this paper \mathbb{K} will denote a field.

Definition 1.1 [19]. Let V denote a vector space over \mathbb{K} with finite positive dimension. By a *Leonard pair* on V we mean an ordered pair A, A^* where $A : V \rightarrow V$ and $A^* : V \rightarrow V$ are linear transformations that satisfy (i) and (ii) below:

- (i) There exists a basis for V with respect to which the matrix representing A is irreducible tridiagonal and the matrix representing A^* is diagonal.
- (ii) There exists a basis for V with respect to which the matrix representing A^* is irreducible tridiagonal and the matrix representing A is diagonal.

Note 1.2. It is a common notational convention to use A^* to represent the conjugate–transpose of A . We are *not* using this convention. In a Leonard pair A, A^* the linear transformations A and A^* are arbitrary subject to (i) and (ii) above.

We refer the reader to [3,9,12–19,21–28,30,31] for background on Leonard pairs. We especially recommend the survey [28]. See [1,2,4–8,10,11,20,29] for related topics.

2. Leonard systems

When working with a Leonard pair, it is convenient to consider a closely related object called a *Leonard system*. To prepare for our definition of a Leonard system, we recall a few concepts from linear algebra. Let d denote a nonnegative integer and let $\text{Mat}_{d+1}(\mathbb{K})$ denote the \mathbb{K} -algebra consisting of all $d+1$ by $d+1$ matrices that have entries in \mathbb{K} . We index the rows and columns by $0, 1, \dots, d$. For the rest of this paper, let \mathcal{A} denote a \mathbb{K} -algebra isomorphic to $\text{Mat}_{d+1}(\mathbb{K})$, and let V denote a simple \mathcal{A} -module. We remark that V is unique up to isomorphism of \mathcal{A} -modules, and that V has dimension $d+1$. Let v_0, v_1, \dots, v_d denote a basis for V . For $X \in \mathcal{A}$ and $Y \in \text{Mat}_{d+1}(\mathbb{K})$, we say Y represents X with respect to v_0, v_1, \dots, v_d whenever $Xv_j = \sum_{i=0}^d Y_{ij}v_i$ for $0 \leq j \leq d$. For $A \in \mathcal{A}$ we say A is *multiplicity-free* whenever it has $d+1$ mutually distinct eigenvalues in \mathbb{K} . Assume A is multiplicity-free. Let $\theta_0, \theta_1, \dots, \theta_d$ denote an ordering of the eigenvalues of A , and for $0 \leq i \leq d$ put

$$E_i = \prod_{\substack{0 \leq j \leq d \\ j \neq i}} \frac{A - \theta_j I}{\theta_i - \theta_j}, \quad (1)$$

where I denotes the identity of \mathcal{A} . We observe (i) $AE_i = \theta_i E_i$ ($0 \leq i \leq d$); (ii) $E_i E_j = \delta_{i,j} E_i$ ($0 \leq i, j \leq d$); (iii) $\sum_{i=0}^d E_i = I$; (iv) $A = \sum_{i=0}^d \theta_i E_i$. Let \mathcal{D} denote the subalgebra of \mathcal{A} generated by A . Using (i)–(iv) we find the sequence E_0, E_1, \dots, E_d is a basis for the \mathbb{K} -vector space \mathcal{D} . We call E_i the *primitive idempotent* of A associated with θ_i . It is helpful to think of these primitive idempotents as follows. Observe

$$V = E_0 V + E_1 V + \dots + E_d V \quad (\text{direct sum}). \quad (2)$$

For $0 \leq i \leq d$, $E_i V$ is the (one dimensional) eigenspace of A in V associated with the eigenvalue θ_i , and E_i acts on V as the projection onto this eigenspace. We remark that the \mathbb{K} -vector space \mathcal{D} has basis $\{A^i | 0 \leq i \leq d\}$ and satisfies $\mathcal{D} = \{X \in \mathcal{A} | AX = XA\}$.

By a *Leonard pair* in \mathcal{A} we mean an ordered pair of elements taken from \mathcal{A} that act on V as a Leonard pair in the sense of Definition 1.1. We now define a Leonard system.

Definition 2.1 [19]. By a *Leonard system* in \mathcal{A} we mean a sequence

$$(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$$

that satisfies (i)–(v) below:

- (i) Each of A, A^* is a multiplicity-free element in \mathcal{A} .
- (ii) E_0, E_1, \dots, E_d is an ordering of the primitive idempotents of A .
- (iii) $E_0^*, E_1^*, \dots, E_d^*$ is an ordering of the primitive idempotents of A^* .
- (iv) For $0 \leq i, j \leq d$,

$$E_i A^* E_j = \begin{cases} 0 & \text{if } |i - j| > 1, \\ \neq 0 & \text{if } |i - j| = 1. \end{cases} \quad (3)$$

- (v) For $0 \leq i, j \leq d$,

$$E_i^* A E_j^* = \begin{cases} 0 & \text{if } |i - j| > 1, \\ \neq 0 & \text{if } |i - j| = 1. \end{cases} \quad (4)$$

Leonard systems are related to Leonard pairs as follows. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Then A, A^* is a Leonard pair in \mathcal{A} [27, Section 3]. Conversely, suppose A, A^* is a Leonard pair in \mathcal{A} . Then each of A, A^* is multiplicity-free [19, Lemma 1.3]. Moreover, there exists an ordering E_0, E_1, \dots, E_d of the primitive idempotents of A , and there exists an ordering $E_0^*, E_1^*, \dots, E_d^*$ of the primitive idempotents of A^* , such that $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ is a Leonard system in \mathcal{A} [27, Lemma 3.3].

3. The space \mathcal{X}

In this paper we consider a subspace of \mathcal{A} defined as follows.

Definition 3.1. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Let \mathcal{X} denote the \mathbb{K} -subspace of \mathcal{A} consisting of the $X \in \mathcal{A}$ such that both

$$E_i X E_j = 0 \quad \text{if } |i - j| > 1, \quad (5)$$

$$E_i^* X E_j^* = 0 \quad \text{if } |i - j| > 1 \quad (6)$$

for $0 \leq i, j \leq d$.

We now state our main result.

Theorem 3.2. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Then the space \mathcal{X} from Definition 3.1 is spanned by

$$I, A, A^*, AA^*, A^*A. \quad (7)$$

Moreover, (7) is a basis for \mathcal{X} provided $d \geq 2$.

The proof of Theorem 3.2 will be given in Section 5.

4. The antiautomorphism \dagger

Associated with a given Leonard system in \mathcal{A} , there is certain antiautomorphism of \mathcal{A} denoted by \dagger and defined below. Recall an *antiautomorphism* of \mathcal{A} is an isomorphism of \mathbb{K} -vector spaces $\sigma : \mathcal{A} \rightarrow \mathcal{A}$ such that $(XY)^\sigma = Y^\sigma X^\sigma$ for all $X, Y \in \mathcal{A}$.

Theorem 4.1 [27, Theorem 7.1]. *Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Then there exists a unique antiautomorphism \dagger of \mathcal{A} such that $A^\dagger = A$ and $A^{*\dagger} = A^*$. Moreover, $X^{\dagger\dagger} = X$ for all $X \in \mathcal{A}$.*

Definition 4.2. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . We let \mathcal{D} denote the subalgebra of \mathcal{A} generated by A . We let \mathcal{D}^* denote the subalgebra of \mathcal{A} generated by A^* .

Lemma 4.3 [28, Lemma 6.3]. *Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} and let \dagger denote the corresponding antiautomorphism of \mathcal{A} from Theorem 4.1. Then referring to Definition 4.2, \dagger fixes everything in \mathcal{D} and everything in \mathcal{D}^* . In particular*

$$E_i^\dagger = E_i, \quad E_i^{*\dagger} = E_i^* \quad (0 \leq i \leq d). \quad (8)$$

5. A basis for \mathcal{X}

In this section we prove Theorem 3.2. We start with a lemma.

Lemma 5.1 [27, Lemma 11.1]. *Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} and let V denote a simple \mathcal{A} -module. Then $E_i V = E_i E_0^* V$ and $E_i^* V = E_i^* E_0 V$ for $0 \leq i \leq d$.*

Corollary 5.2. *Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Then for $Y \in \mathcal{A}$ the following hold for $0 \leq i \leq d$:*

- (i) $Y E_i = 0$ if and only if $Y E_i E_0^* = 0$.
- (ii) $Y E_i^* = 0$ if and only if $Y E_i^* E_0 = 0$.

Corollary 5.3. *Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Then for $Y \in \mathcal{A}$ the following hold for $0 \leq i \leq d$:*

- (i) $E_i Y = 0$ if and only if $E_0^* E_i Y = 0$.
- (ii) $E_i^* Y = 0$ if and only if $E_0 E_i^* Y = 0$.

Proof. Apply \dagger to the equations in Corollary 5.2, and use Lemma 4.3. \square

Definition 5.4. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . For $0 \leq i \leq d$ we let θ_i (respectively θ_i^*) denote the eigenvalue of A (respectively A^*) associated with E_i (respectively E_i^*). We note that the scalars $\theta_0, \theta_1, \dots, \theta_d$ (respectively $\theta_0^*, \theta_1^*, \dots, \theta_d^*$) are mutually distinct and contained in \mathbb{K} .

Proposition 5.5. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} and let \mathcal{X} denote the subspace of \mathcal{A} from Definition 3.1. Let $X \in \mathcal{X}$ such that $XE_0^* = 0$ and $XAE_0^* = 0$. Then $X = 0$.

Proof. First assume $d = 0$. Then $E_0^* = I$ and the result follows. For the rest of this proof assume $d \geq 1$. We assume $X \neq 0$ and get a contradiction.

In the equation $I = \sum_{i=0}^d E_i^*$ we multiply each term on the right by AE_0^* and simplify the result using (4) to obtain $AE_0^* = E_0^*AE_0^* + E_1^*AE_0^*$; expanding $XAE_0^* = 0$ using this and $XE_0^* = 0$ we find $XE_1^*AE_0^* = 0$. Let V denote a simple \mathcal{A} -module and observe $XE_1^*AE_0^*V = 0$. Note that $E_1^*V = E_1^*AE_0^*V$, since $E_1^*AE_0^*V \subseteq E_1^*V$, $\dim E_1^*V = 1$, and $E_1^*AE_0^*V \neq 0$ in view of (4). By the above comments $XE_1^*V = 0$ so $XE_1^* = 0$. In the equation $I = \sum_{i=0}^d E_i^*$ we multiply each term on the left by E_0^*X and simplify the result using (6) to find $E_0^*X = E_0^*XE_0^* + E_0^*XE_1^*$; now $E_0^*X = 0$ since each of XE_0^* , XE_1^* is zero.

Since $X \neq 0$ there exist integers i, j ($0 \leq i, j \leq d$) such that $E_iXE_j \neq 0$. Define

$$r = \min\{\min\{i, j\} | 0 \leq i, j \leq d, E_iXE_j \neq 0\}.$$

First assume $r = d$, so that $E_dXE_d \neq 0$ and each of E_iXE_d, E_dXE_i is zero for $0 \leq i \leq d-1$. In the equation $I = \sum_{i=0}^d E_i$ we multiply each term on the left by E_dX and simplify to get $E_dX = E_dXE_d$. By this and since $XE_0^* = 0$ we find $E_dXE_dE_0^* = 0$. Now $E_dXE_d = 0$ by Corollary 5.2(i), for a contradiction.

Next assume $r \leq d-1$. Note that for $0 \leq i \leq r-1$ we have $E_rXE_i = 0$ and $E_iXE_r = 0$. We now show that each of E_rXE_r and E_rXE_{r+1} is zero. In the equation $I = \sum_{i=0}^d E_i$ we multiply each term on the left by E_rX . We simplify the result using (5) and our above comments to find

$$E_rX = E_rXE_r + E_rXE_{r+1}. \quad (9)$$

In this equation we multiply each term on the right by E_0^* and use $XE_0^* = 0$ to find

$$E_rXE_rE_0^* + E_rXE_{r+1}E_0^* = 0. \quad (10)$$

We multiply each term of (9) on the right by A and use $E_iA = \theta_iE_i$ ($0 \leq i \leq d$) to find $E_rXA = \theta_rE_rXE_r + \theta_{r+1}E_rXE_{r+1}$. In this equation we multiply each term on the right by E_0^* and use $XAE_0^* = 0$ to find

$$\theta_rE_rXE_rE_0^* + \theta_{r+1}E_rXE_{r+1}E_0^* = 0. \quad (11)$$

Solving the linear system (10) and (11), we find $E_rXE_rE_0^* = 0$ and $E_rXE_{r+1}E_0^* = 0$. By this and Corollary 5.2(i) we find $E_rXE_r = 0$ and $E_rXE_{r+1} = 0$. Next we show $E_{r+1}XE_r = 0$. We mentioned earlier that $E_iXE_r = 0$ for $0 \leq i \leq r-1$. In the equation $I = \sum_{i=0}^d E_i$ we multiply each term on the right by XE_r . We simplify the result using (5) and our above comments to find $XE_r = E_{r+1}XE_r$. In this equation we multiply each term on the left by E_0^* and use $E_0^*X = 0$ to find $E_0^*E_{r+1}XE_r = 0$, so $E_{r+1}XE_r = 0$ in view of Corollary 5.3(i). We have now shown that each of $E_rXE_r, E_rXE_{r+1}, E_{r+1}XE_r$ is zero, contradicting the definition of r . We conclude $X = 0$. \square

Corollary 5.6. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Then the space \mathcal{X} from Definition 3.1 has dimension at most 5.

Proof. We assume $d \geq 2$; otherwise $\dim \mathcal{A} \leq 4$ and the result follows. We define linear maps $\pi_0 : \mathcal{X} \rightarrow \mathcal{X}E_0^*$ and $\pi_1 : \mathcal{X} \rightarrow \mathcal{X}AE_0^*$ by

$$\pi_0(X) = XE_0^*, \quad \pi_1(X) = XAE_0^* \quad (X \in \mathcal{X}).$$

For $i = 0, 1$ let K_i denote the kernel of π_i . We compute the dimensions of K_0 and K_1 . First observe

$$\dim E_i^* \mathcal{A} E_j^* = 1 \quad (0 \leq i, j \leq d).$$

We have $\mathcal{X}E_0^* = E_0^* \mathcal{X} E_0^* + E_1^* \mathcal{X} E_0^*$ in view of (6); therefore $\dim \mathcal{X}E_0^* \leq 2$ so

$$\dim K_0 \geq \dim \mathcal{X} - 2. \quad (12)$$

Combining (4) and (6) we routinely obtain

$$\mathcal{X}AE_0^* \subseteq E_0^* \mathcal{A} E_0^* + E_1^* \mathcal{A} E_0^* + E_2^* \mathcal{A} E_0^*.$$

Therefore, $\dim \mathcal{X}AE_0^* \leq 3$ so

$$\dim K_1 \geq \dim \mathcal{X} - 3. \quad (13)$$

The intersection of K_0 and K_1 is zero by Proposition 5.5; therefore

$$\dim K_0 + \dim K_1 \leq \dim \mathcal{X}. \quad (14)$$

Combining (12)–(14) we find $\dim \mathcal{X} \leq 5$ as desired. \square

Proof of Theorem 3.2. Comparing (3), (4) and (5), (6) we see that each of the elements (7) is contained in \mathcal{X} . We must show they actually span \mathcal{X} , and that they are linearly independent provided $d \geq 2$. First assume $d = 0$. Then the assertion is obvious. Next assume $d = 1$. Then one routinely verifies that $\mathcal{X} = \mathcal{A}$ is spanned by the elements (7). Finally assume $d \geq 2$. In view of Corollary 5.6, it suffices to show that the elements (7) are linearly independent. Suppose

$$eI + fA + f^*A^* + gAA^* + g^*A^*A = 0 \quad (15)$$

for some scalars e, f, f^*, g, g^* in \mathbb{K} . We show each of e, f, f^*, g, g^* is zero. For $1 \leq i \leq d$ we multiply each term in (15) on the left by E_{i-1}^* and the right by E_i^* to obtain

$$(f + g\theta_i^* + g^*\theta_{i-1}^*)E_{i-1}^*AE_i^* = 0.$$

By this and since $E_{i-1}^*AE_i^*$ is nonzero we find

$$f + g\theta_i^* + g^*\theta_{i-1}^* = 0 \quad (1 \leq i \leq d). \quad (16)$$

For $1 \leq i \leq d$ we multiply each term in (15) on the left by E_i^* and the right by E_{i-1}^* to obtain

$$(f + g\theta_{i-1}^* + g^*\theta_i^*)E_i^*AE_{i-1}^* = 0.$$

By this and since $E_i^*AE_{i-1}^*$ is nonzero we find

$$f + g\theta_{i-1}^* + g^*\theta_i^* = 0 \quad (1 \leq i \leq d). \quad (17)$$

Combining (16) at $i = 1$ and (17) at $i = 1, 2$ we routinely find that each of f, g, g^* is zero. Interchanging the roles of A and A^* in the above argument we find $f^* = 0$. Now (15) becomes $eI = 0$ so $e = 0$. We have now shown that each of e, f, f^*, g, g^* is zero, and the result follows. \square

6. The linear maps Υ and Υ^*

In this section we discuss some linear maps $\Upsilon : \mathcal{X} \rightarrow \mathcal{D}$ and $\Upsilon^* : \mathcal{X} \rightarrow \mathcal{D}^*$ that we find attractive. To motivate things we recall some results by the second author and Vidunas.

Lemma 6.1 [30, Theorem 1.5]. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Then there exists a sequence of scalars $\beta, \gamma, \gamma^*, \varrho, \varrho^*, \omega, \eta, \eta^*$ taken from \mathbb{K} such that both

$$A^2 A^* - \beta A A^* A + A^* A^2 - \gamma(AA^* + A^*A) - \varrho A^* = \gamma^* A^2 + \omega A + \eta I, \quad (18)$$

$$A^* A^2 - \beta A^* A A^* + A A^{*2} - \gamma^*(A^*A + A A^*) - \varrho^* A = \gamma A^{*2} + \omega A^* + \eta^* I. \quad (19)$$

Moreover, the sequence is uniquely determined by the Leonard system provided $d \geq 3$.

Note 6.2. Eqs. (18) and (19) first appeared in [32]; they are called the Askey–Wilson relations.

We have a comment.

Lemma 6.3 [30, Theorem 4.5]. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Then referring to Definition 5.4 and Lemma 6.1 we have

$$\beta + 1 = \frac{\theta_{i-2} - \theta_{i+1}}{\theta_{i-1} - \theta_i} = \frac{\theta_{i-2}^* - \theta_{i+1}^*}{\theta_{i-1}^* - \theta_i^*} \quad (2 \leq i \leq d-1), \quad (20)$$

$$\gamma = \theta_{i-1} - \beta\theta_i + \theta_{i+1} \quad (1 \leq i \leq d-1), \quad (21)$$

$$\gamma^* = \theta_{i-1}^* - \beta\theta_i^* + \theta_{i+1}^* \quad (1 \leq i \leq d-1), \quad (22)$$

$$\varrho = \theta_{i-1}^2 - \beta\theta_{i-1}\theta_i + \theta_i^2 - \gamma(\theta_{i-1} + \theta_i) \quad (1 \leq i \leq d), \quad (23)$$

$$\varrho^* = \theta_{i-1}^{*2} - \beta\theta_{i-1}^*\theta_i^* + \theta_i^{*2} - \gamma^*(\theta_{i-1}^* + \theta_i^*) \quad (1 \leq i \leq d). \quad (24)$$

Theorem 6.4. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Let the spaces \mathcal{X} and \mathcal{D} be as in Definitions 3.1 and 4.2, respectively. Then there exists a \mathbb{K} -linear map $\Upsilon : \mathcal{X} \rightarrow \mathcal{D}$ that satisfies

$$\Upsilon(X) = A^2 X - \beta A X A + X A^2 - \gamma(A X + X A) - \varrho X \quad (25)$$

for all $X \in \mathcal{X}$. Moreover,

$$\Upsilon(I) = (2 - \beta)A^2 - 2\gamma A - \varrho I, \quad (26)$$

$$\Upsilon(A) = (2 - \beta)A^3 - 2\gamma A^2 - \varrho A, \quad (27)$$

$$\Upsilon(A^*) = \gamma^* A^2 + \omega A + \eta I, \quad (28)$$

$$\Upsilon(AA^*) = \gamma^* A^3 + \omega A^2 + \eta A, \quad (29)$$

$$\Upsilon(A^*A) = \gamma^* A^3 + \omega A^2 + \eta A. \quad (30)$$

Proof. Certainly, there exists a \mathbb{K} -linear map $\Upsilon : \mathcal{X} \rightarrow \mathcal{A}$ that satisfies (25). Using (18) we find Υ satisfies (26)–(30). Combining (26)–(30) and Theorem 3.2 we find $\Upsilon(X) \in \mathcal{D}$ for all $X \in \mathcal{X}$, and the result follows. \square

Interchanging the roles of A and A^* in Theorem 6.4 we obtain:

Theorem 6.5. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Let the spaces \mathcal{X} and \mathcal{D}^* be as in Definitions 3.1 and 4.2, respectively. Then there exists a \mathbb{K} -linear map $\Upsilon^* : \mathcal{X} \rightarrow \mathcal{D}^*$ that satisfies

$$\Upsilon^*(X) = A^{*2}X - \beta A^*XA^* + XA^{*2} - \gamma^*(A^*X + XA^*) - \varrho^*X \quad (31)$$

for all $X \in \mathcal{X}$. Moreover,

$$\begin{aligned} \Upsilon^*(I) &= (2 - \beta)A^{*2} - 2\gamma^*A^* - \varrho^*I, \\ \Upsilon^*(A^*) &= (2 - \beta)A^{*3} - 2\gamma^*A^{*2} - \varrho^*A^*, \\ \Upsilon^*(A) &= \gamma A^{*2} + \omega A^* + \eta^*I, \\ \Upsilon^*(A^*A) &= \gamma A^{*3} + \omega A^{*2} + \eta^*A^*, \\ \Upsilon^*(AA^*) &= \gamma A^{*3} + \omega A^{*2} + \eta^*A^*. \end{aligned}$$

We have a comment concerning the image and kernel of Υ .

Lemma 6.6. Referring to Theorem 6.4 the following (i)–(iii) hold.

- (i) $\text{Span}\{AA^* - A^*A\} \subseteq \text{Ker}(\Upsilon)$.
- (ii) $\text{Im}(\Upsilon) \subseteq \text{Span}\{I, A, A^2, A^3\}$.
- (iii) Assume $d \geq 3$. Then equality holds in (i) if and only if equality holds in (ii).

Proof. (i) and (ii): Immediate from Theorem 6.4.

(iii): Use Theorem 3.2 and elementary linear algebra. \square

Interchanging the roles of A and A^* in Lemma 6.6 we obtain:

Lemma 6.7. Referring to Theorem 6.5 the following (i)–(iii) hold.

- (i) $\text{Span}\{AA^* - A^*A\} \subseteq \text{Ker}(\Upsilon^*)$.
- (ii) $\text{Im}(\Upsilon^*) \subseteq \text{Span}\{I, A^*, A^{*2}, A^{*3}\}$.
- (iii) Assume $d \geq 3$. Then equality holds in (i) if and only if equality holds in (ii).

Referring to Lemmas 6.6 and 6.7 it appears that we have equality in (i) and (ii) for most Leonard systems but not all. Below we give an example where equality is not attained.

Definition 6.8. Let $(A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . We say this Leonard system is *bipartite* (respectively *dual bipartite*) whenever $E_i^*AE_i^* = 0$ (respectively $E_iA^*E_i = 0$) for $0 \leq i \leq d$.

Lemma 6.9. Let $\Phi = (A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system in \mathcal{A} . Then referring to Theorems 6.4, 6.5 and Definition 6.8 the following (i), (ii) hold provided $d \geq 3$.

- (i) Assume Φ is bipartite. Then

$$\begin{aligned} \text{Ker}(\Upsilon^*) &= \text{Span}\{A, AA^*, A^*A\}, \\ \text{Im}(\Upsilon^*) &= \text{Span}\{B^*, A^*B^*\}, \end{aligned}$$

where $B^* = (2 - \beta)A^{*2} - 2\gamma^*A^* - \varrho^*I$.

(ii) Assume Φ is dual bipartite. Then

$$\begin{aligned}\text{Ker}(\Upsilon) &= \text{Span}\{A^*, A^*A, AA^*\}, \\ \text{Im}(\Upsilon) &= \text{Span}\{B, AB\},\end{aligned}$$

$$\text{where } B = (2 - \beta)A^2 - 2\gamma A - \varrho I.$$

Proof. (ii): By [12] and [30, Theorem 5.3] each of γ^* , ω , η is zero. By this and Theorem 6.4 we have $\text{Ker}(\Upsilon) \supseteq \text{Span}\{A^*, A^*A, AA^*\}$ and $\text{Im}(\Upsilon) = \text{Span}\{B, AB\}$. To show $\text{Ker}(\Upsilon) = \text{Span}\{A^*, A^*A, AA^*\}$ it suffices to show that B and AB are linearly independent. Suppose B and AB are linearly dependent. Then $B = 0$ since the elements I, A, A^2, A^3 are linearly independent. Since $d \geq 3$ there exists an integer i such that $1 \leq i \leq d - 1$. Multiplying each term in the equation $B = (2 - \beta)A^2 - 2\gamma A - \varrho I$ by E_i and simplifying we find E_i times

$$(2 - \beta)\theta_i^2 - 2\gamma\theta_i - \varrho \tag{32}$$

is zero. Of course E_i is not zero so (32) is zero. Using (21) and (23) we routinely find (32) is equal to

$$(\theta_i - \theta_{i-1})(\theta_i - \theta_{i+1})$$

and is therefore nonzero. This is a contradiction and the result follows. \square

Open problem: Referring to Lemmas 6.6 and 6.7, precisely determine the set of Leonard systems for which equality holds in (i) and (ii).

References

- [1] H. Alnajjar, B. Curtin, A family of tridiagonal pairs, *Linear Algebra Appl.* 390 (2004) 369–384.
- [2] H. Alnajjar, B. Curtin, A family of tridiagonal pairs related to the quantum affine algebra $U_q(\widehat{sl}_2)$, *Electron. J. Linear Algebra* 13 (2005) 1–9.
- [3] B. Hartwig, Three mutually adjacent Leonard pairs, *Linear Algebra Appl.* 408 (2005) 19–39, Available from: [arXiv:math.AC/0508415](https://arxiv.org/abs/math.AC/0508415).
- [4] T. Ito, K. Tanabe, P. Terwilliger, Some algebra related to P - and Q -polynomial association schemes, *Codes and Association Schemes* (Piscataway, NJ, 1999), American Mathematical Society, Providence, RI, 2001, pp. 167–192, Available from: [arXiv:math.CO/0406556](https://arxiv.org/abs/math.CO/0406556).
- [5] T. Ito, P. Terwilliger, The shape of a tridiagonal pair, *J. Pure Appl. Algebra* 188 (2004) 145–160, Available from: [arXiv:math.QA/0304244](https://arxiv.org/abs/math.QA/0304244).
- [6] T. Ito, P. Terwilliger, Tridiagonal pairs and the quantum affine algebra $U_q(\widehat{sl}_2)$, *Ramanujan J.*, in press. Available from: [arXiv:math.QA/0310042](https://arxiv.org/abs/math.QA/0310042).
- [7] T. Ito, P. Terwilliger, Two non-nilpotent linear transformations that satisfy the cubic q -Serre relations, *J. Algebra Appl.*, submitted for publication. Available from: [arXiv:math.QA/0508398](https://arxiv.org/abs/math.QA/0508398).
- [8] T. Ito, P. Terwilliger, C. Weng, The quantum algebra $U_q(sl_2)$ and its equitable presentation, *J. Algebra* 298 (2006) 284–301, Available from: [arXiv:math.QA/0507477](https://arxiv.org/abs/math.QA/0507477).
- [9] K. Nomura, Tridiagonal pairs and the Askey–Wilson relations, *Linear Algebra Appl.* 397 (2005) 99–106.
- [10] K. Nomura, A refinement of the split decomposition of a tridiagonal pair, *Linear Algebra Appl.* 403 (2005) 1–23.
- [11] K. Nomura, Tridiagonal pairs of height one, *Linear Algebra Appl.* 403 (2005) 118–142.
- [12] K. Nomura, P. Terwilliger, Balanced Leonard pairs, *Linear Algebra Appl.*, in press, doi:10.1016/j.laa.2006.06.025, Available from: [arXiv:math.RA/0506219](https://arxiv.org/abs/math.RA/0506219).
- [13] K. Nomura, P. Terwilliger, Some trace formulae involving the split sequences of a Leonard pair, *Linear Algebra Appl.* 413 (2006) 189–201, Available from: [arXiv:math.RA/0508407](https://arxiv.org/abs/math.RA/0508407).
- [14] K. Nomura, P. Terwilliger, The determinant of $AA^* - A^*A$ for a Leonard pair A, A^* , *Linear Algebra Appl.* 416 (2006) 880–889, Available from: [arXiv:math.RA/0511641](https://arxiv.org/abs/math.RA/0511641).

- [15] K. Nomura, P. Terwilliger, Matrix units associated with the split basis of a Leonard pair, *Linear Algebra Appl.*, in press, doi:10.1016/j.laa.2006.03.009, Available from: <arXiv:math.RA/0602416>.
- [16] A.A. Pascasio, On the multiplicities of the primitive idempotents of a Q -polynomial distance-regular graph, *Eur. J. Combin.* 23 (2002) 1073–1078.
- [17] P. Terwilliger, The subconstituent algebra of an association scheme I, *J. Algebraic Combin.* 1 (1992) 363–388.
- [18] P. Terwilliger, The subconstituent algebra of an association scheme III, *J. Algebraic Combin.* 2 (1993) 177–210.
- [19] P. Terwilliger, Two linear transformations each tridiagonal with respect to an eigenbasis of the other, *Linear Algebra Appl.* 330 (2001) 149–203, Available from: <arXiv:math.RA/0406555>.
- [20] P. Terwilliger, Two relations that generalize the q -Serre relations and the Dolan–Grady relations, *Physics and Combinatorics 1999* (Nagoya), World Scientific Publishing, River Edge, NJ, 2001, pp. 377–398, Available from: <arXiv:math.QA/0307016>.
- [21] P. Terwilliger, Leonard pairs from 24 points of view, *Rocky Mountain J. Math.* 32 (2) (2002) 827–888, Available from: <arXiv:math.RA/0406577>.
- [22] P. Terwilliger, Two linear transformations each tridiagonal with respect to an eigenbasis of the other; the T - D and the LB - UB canonical form, *J. Algebra* 291 (2005) 1–45, Available from: <arXiv:math.RA/0304077>.
- [23] P. Terwilliger, Introduction to Leonard pairs, *J. Comput. Appl. Math.* (2) (2003) 463–475.
- [24] P. Terwilliger, Introduction to Leonard pairs and Leonard systems, *Sūrikaiseki-kenkyūsho Kōkyūroku* 1109 (1999) 67–79, *Algebraic Combinatorics* (Kyoto, 1999).
- [25] P. Terwilliger, Two linear transformations each tridiagonal with respect to an eigenbasis of the other; comments on the split decomposition, *J. Comput. Appl. Math.* 178 (2005) 437–452, Available from: <arXiv:math.RA/0306290>.
- [26] P. Terwilliger, Two linear transformations each tridiagonal with respect to an eigenbasis of the other; comments on the parameter array, *Des. Codes Cryptogr.* 34 (2005) 307–332, Available from: <arXiv:math.RA/0306291>.
- [27] P. Terwilliger, Leonard pairs and the q -Racah polynomials, *Linear Algebra Appl.* 387 (2004) 235–276, Available from: <arXiv:math.QA/0306301>.
- [28] P. Terwilliger, Two linear transformations each tridiagonal with respect to an eigenbasis of the other: an algebraic approach to the Askey scheme of orthogonal polynomials, *Lecture notes for the Summer School on Orthogonal Polynomials and Special Functions*, Universidad Carlos III de Madrid, Leganes, Spain, July 8–July 18, 2004. Available from: <arXiv:math.QA/0408390>.
- [29] P. Terwilliger, The equitable presentation for the quantum group $U_q(g)$ associated with a symmetrizable Kac–Moody algebra g , *J. Algebra* 298 (2006) 302–319, Available from: <arXiv:math.QA/0507478>.
- [30] P. Terwilliger, R. Vidunas, Leonard pairs and the Askey–Wilson relations, *J. Algebra Appl.* 3 (2004) 411–426, Available from: <arXiv:math.QA/0305356>.
- [31] R. Vidunas, Normalized Leonard pairs and Askey–Wilson relations. Available from: <arXiv:math.RA/0505041>.
- [32] A.S. Zhedanov, “Hidden symmetry” of Askey–Wilson polynomials, *Teoret. Mat. Fiz.* 89 (1991) 190–204.